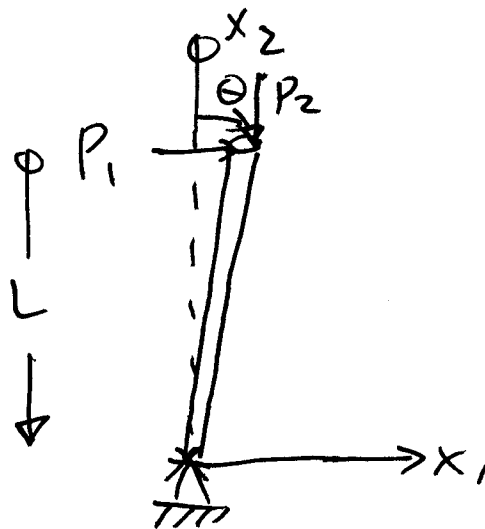


Unified Engineering Problem Set
Week #10 Spring, 2008

SOLUTIONS

110.1 Begin with general configuration:



(a) The rod can deform along its length proportional to its stiffness

$$\Delta L = P_2 / k_L$$

with $k_L = EA/L$ as defined

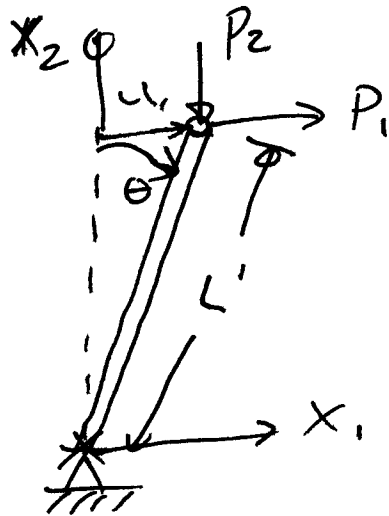
Thus the length, L' , at any time is:

$$L' = L - \Delta L \Rightarrow$$

$$L' = L - \frac{P_2}{EA/L} = L \left(1 - \frac{P_2}{EA}\right)$$

This comes with positive ΔL and P_2 defined as being in negative x_2 direction

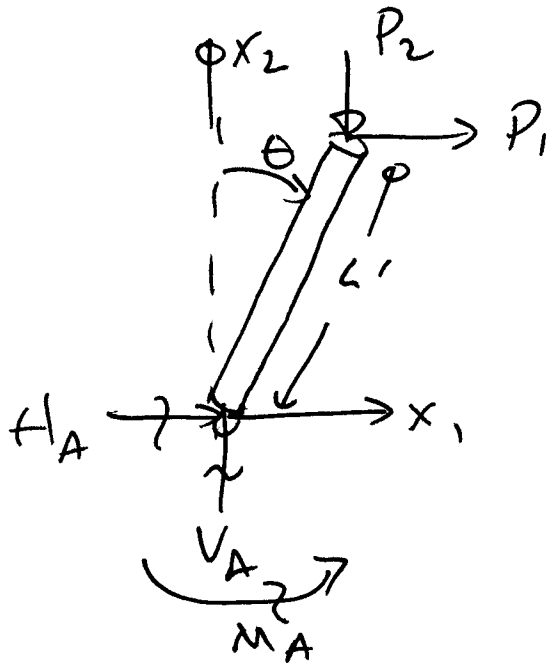
So this "redraws" the configuration as:



The distance u_1 is $L' \sin \theta$
and again assuming small deformations
 $\Rightarrow u_1 = L' \theta$

The distance in the x_2 - direction is $L' \cos \theta$
 $\Rightarrow \dots = L'$

Now draw the Free Body Diagram and take equilibrium:



$$\Sigma F_{x_1} = 0 \xrightarrow{+} \Rightarrow P_1 + H_A = 0 \Rightarrow H_A = -P_1$$

$$\Sigma F_{x_2} = 0 \uparrow \Rightarrow V_A - P_2 = 0 \Rightarrow V_A = +P_2$$

$$\Sigma M_A = 0 \left(\begin{array}{l} \uparrow \\ \text{sign} \end{array} \right) \Rightarrow -P_1 L' - P_2 L' \theta + k_T \theta = 0$$

acts in opposite
direction from $+\theta$
 $+\theta$ is CW \Rightarrow moment is
CW $\Rightarrow +$

using the latter equation:

$$\theta (k_T - P_2 L') = P_1 L'$$

and thus:

$$\theta \left(\frac{k_T - P_2 L'}{L'} \right) = P_1$$

giving

$$\theta \left(\frac{k_T}{L'} - P_2 \right) = P_1$$

use the expression for L' to get:

$$\Theta \left[\frac{k_T}{L \left(1 - \frac{P_2}{EA}\right)} - P_2 \right] = P_1$$

Basic equation governing rotational angle, Θ

(b) This expression is basically the form for the effective stiffness, k_{eff} :

$$\Theta k_{eff} = P_1$$

$$\Rightarrow k_{eff} = \frac{k_T}{L \left(1 - \frac{P_2}{EA}\right)} - P_2$$

and when $k_{eff} = 0$, instability occurs so the basic equation is:

$$\frac{k_T}{L \left(1 - \frac{P_2}{EA}\right)} - P_2 = 0 \quad (*)$$

working this further:

$$\frac{1}{L \left(1 - \frac{P_2}{EA}\right)} \left[k_T - P_2 L + P_2^2 \frac{L}{EA} \right] = 0$$

Within the brackets is a quadratic in the applied axial load, P_2 . Solve this to find the critical value of P_2 for the instability: P_{2cr}

first multiply the expressions by $EA/L/EA/L$:

$$\Rightarrow \frac{1}{EA - P_2} \left[P_2^2 - P_2 EA + k_T \frac{EA}{L} \right] = 0$$

solve the quadratic to be zero via:

$$P_{2cr} = \frac{-b \pm \sqrt{b^2 - 4ac}}{2a}$$

with: $a = 1$
 $b = -EA$
 $c = k_T \frac{EA}{L}$

$$\Rightarrow P_{2cr} = \frac{EA \pm \sqrt{(EA)^2 - 4k_T \frac{EA}{L}}}{2}$$

expression for critical value of P_2

(Note: can not have $P_2 = EA$ since the term $\frac{1}{EA - P_2}$ would approach ∞)

\Rightarrow that would need to be considered

(c) Now use the ratio:

$$n = EA / (k_T L)$$

$$\Rightarrow EA = n \frac{k_T}{L}$$

use this in the expression for P_{2cr} :

$$P_{2cr} = \frac{n \frac{k_T}{L} \pm \sqrt{n^2 \left(\frac{k_T}{L}\right)^2 - 4n \left(\frac{k_T}{L}\right)^2}}{2}$$

$$\Rightarrow P_{2cr} = \frac{k_T}{2L} (n \pm \sqrt{n^2 - 4n})$$

Now it can be seen that to get a real value, must have:

$$n^2 - 4n \geq 0$$

$$\Rightarrow n(n-4) \geq 0$$

$$\Rightarrow n \geq 4$$

→ Behavior: below this value, one gets collapse as the rod is not stiff enough

→ Now consider limits on the values...

for $n=4$:

$$P_{2cr} = \frac{k_T}{2L} (4) = 2 \frac{k_T}{L}$$

→ Behavior: twice the value for the case treated in lecture of infinite stiffness

for $n = \infty$

→ There are two behaviors depending upon how ∞ is approached

Case ① For the case of $-\sqrt{\quad}$, one gets a look at $EA \rightarrow \infty$:

$$P_{2cr} = \frac{k_T}{2L} (n - \sqrt{n^2 - 4n})$$

for this case one must look at the residual from $\sqrt{n^2 - 4n}$ and this is $(n-2)$.

So this becomes:

$$P_{2cr} = \frac{k_T}{2L} (2) = \frac{k_T}{L}$$

→ Behavior: Same as value derived in lecture. One can also look at equation (*) and see that if $EA \rightarrow \infty$,

$$\text{then } \frac{P_2}{EA} \rightarrow 0 \quad \text{so } \left. \begin{array}{l} \frac{k_T}{L} - P_2 = 0 \\ \Rightarrow P_2 = \frac{k_T}{L} \end{array} \right\} \begin{array}{l} \text{as in} \\ \text{lecture} \end{array}$$

Case ② For the case of $+\sqrt{\quad}$, look at

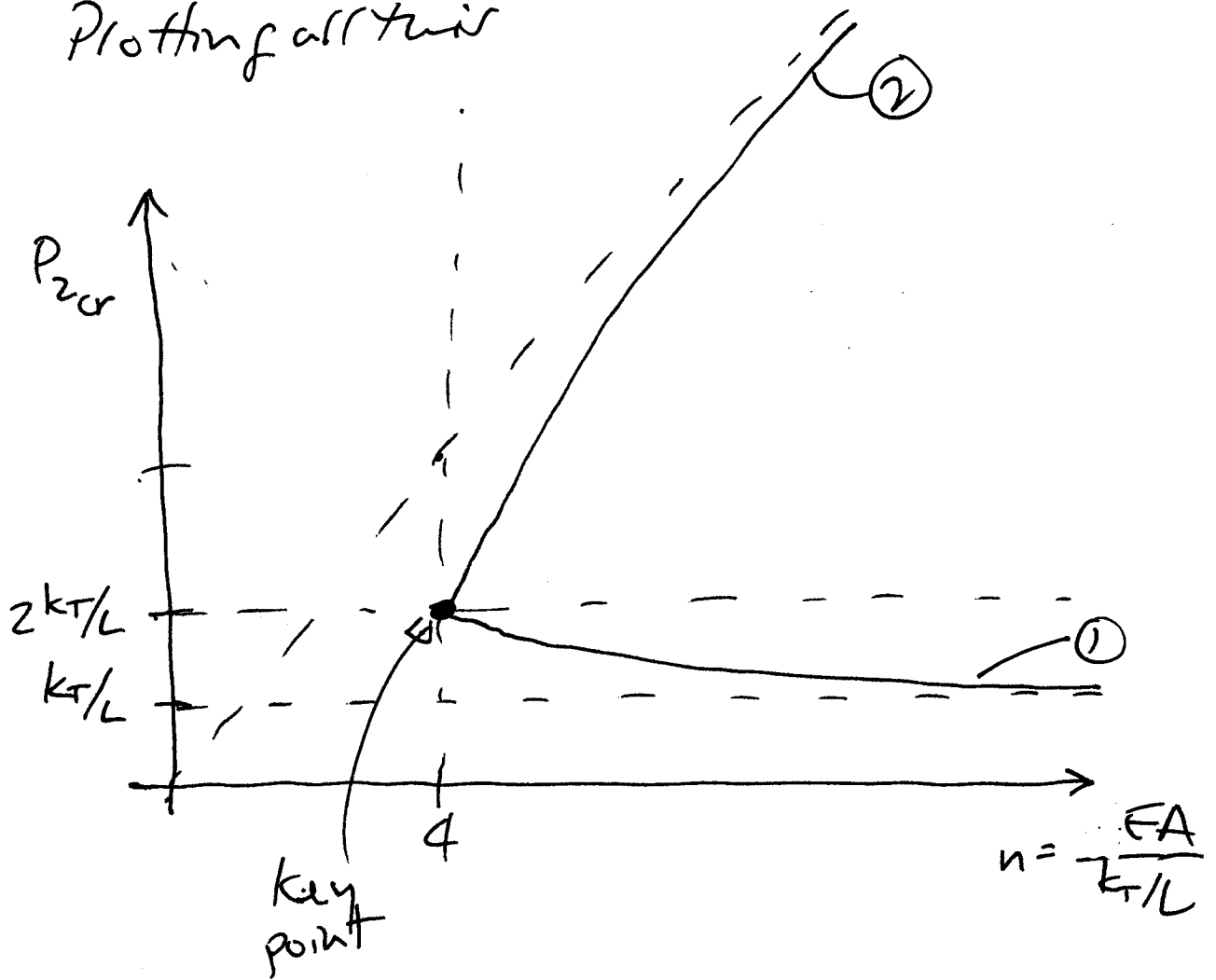
$$\frac{AE}{k_T/L} \rightarrow \infty$$

$$\text{Now: } P = \frac{kT}{2r} (n + \sqrt{n^2})$$

$$\Rightarrow P_{2cr} = \frac{kT}{L} n$$

→ Behavior: An asymptotic line proportional to kT/L

Plotting all this



And always

$$P_{cr} > \frac{kT}{L}$$

$$P_{cr} < EA$$

M 10.2

General configuration:



use a general structural property characterization:

Cross-section Area = A

Moment of Inertia = I

modulus = E

end Load = P (related to pressure and Area)

This allows one to characterize the behavior independent of the specifics of the cross-section. This can be probed after determining the expression for P_{cr} , etc. in terms of A , I , and E by expressing A , I , and E in terms of the details.

Start with the basic governing equation:

$$\frac{d^2 u_3}{dx_1^2} + \frac{P}{EI} u_3 = 0 \quad (1)$$

with the general homogeneous solution:

$$u_3 = A \sin\left(\sqrt{\frac{P}{EI}} x_1\right) + B \left(\cos\sqrt{\frac{P}{EI}} x_1\right) + C + D x_1 \quad (2)$$

Now look at the boundary conditions.

• At the clamped end ($x_1 = 0$): $u_3 = 0$
 $\frac{du_3}{dx_1} = 0$

• At the free end with applied load ($x_1 = L$):

[find from the notes]

$$M = 0 \Rightarrow \frac{d^2 u_3}{dx_1^2} = 0$$

$$S = -P \frac{du_3}{dx_1} \text{ and } S = EI \frac{d^3 u_3}{dx_1^3}$$

So the condition is:

$$\frac{d^3 u_3}{dx_1^3} = -\frac{P}{EI} \frac{du_3}{dx_1}$$

→ To facilitate writing the solution, let us represent:

$$\lambda = \sqrt{\frac{P}{EI}}$$

(Note: A bit different than original $\lambda^2 = -\frac{P}{EI}$)

So (1) becomes:

$$\frac{d^2 u_3}{dx_1^2} + \lambda^2 u_3 = 0 \quad (1)$$

and (2) becomes:

$$u_3 = A \sin \lambda x_1 + B \cos \lambda x_1 + C + Dx_1 \quad (2)$$

To use the boundary conditions in the basic solution, need derivatives of (2). So:

$$\frac{du_3}{dx_1} = \lambda A \cos \lambda x_1 - \lambda B \sin \lambda x_1 + D$$

$$\frac{d^2 u_3}{dx_1^2} = -\lambda^2 A \sin \lambda x_1 - \lambda^2 B \cos \lambda x_1$$

$$\frac{d^3 u_3}{dx_1^3} = -\lambda^3 A \cos \lambda x_1 + \lambda^3 B \sin \lambda x_1$$

Now apply each of the 4 Boundary Conditions to get 4 equations:

$$\textcircled{a} \quad x_1 = 0, \quad u_3 = 0 \Rightarrow B + C = 0 \quad (3)$$

$$\textcircled{a} \quad x_1 = 0, \quad \frac{du_3}{dx_1} = 0 \Rightarrow \lambda A + D = 0 \quad (4)$$

$$\textcircled{a} \quad x_1 = L, \quad \frac{d^3 u_3}{dx_1^3} = -\lambda^2 \frac{du_3}{dx_1}$$

as expressed

$$\begin{aligned} \Rightarrow -\lambda^3 A \cos \lambda L + \lambda^3 B \sin \lambda L \\ = -\lambda^3 A \cos \lambda L + \lambda^3 B \sin \lambda L + \lambda^2 D \quad (5) \end{aligned}$$

$$\Rightarrow \boxed{D = 0} \quad (5)$$

and @ $x_1 = L$, $\frac{d^2 u_3}{dx_1^2} = 0$

$$\Rightarrow -\lambda^2 A \sin \lambda L - \lambda^2 B \cos \lambda L = 0$$

$$\text{giving: } A \sin \lambda L + B \cos \lambda L = 0 \quad (6)$$

→ Now manipulate these equations:

from (4): $\lambda A = -D$

$$\Rightarrow \lambda A = 0$$

$$\Rightarrow \boxed{A = 0}$$

from (6):

$$B \cos \lambda L = 0$$

This gives 2 possible solutions:

$$B = 0 \quad (\text{trivial} \dots u_3 = 0)$$

or

$$\cos\left(\sqrt{\frac{P}{EI}} L\right) = 0$$

$$\Rightarrow \sqrt{\frac{P}{EI}} L = \frac{n\pi}{2} \quad (\text{for } n \text{ odd})$$

$$\text{So: } \boxed{P = \frac{n^2 \pi^2}{4L^2} EI} \quad (\text{for } n \text{ odd})$$

The lowest load occurs for $n=1$

$$\boxed{P_{cr} = \frac{\pi^2 EI}{4L^2}} \quad \text{buckling load}$$

In general, the form for P_{cr} is:

$$P_{cr} = c \frac{\pi^2 EI}{L^2}$$

where c = coefficient of (edge) fixity

So, in this case: $c = 1/4$

To get the mode, go to equation (3) fixing:

$$B + C = 0$$

$$\Rightarrow B = -C$$

Use this in the expression for u_3 of equation (2), along with the other results:

$$\text{gives: } u_3 = B \left[\cos \left(\sqrt{\frac{n^2 \pi^2 EI / 4L^2}{EI}} x_1 \right) - 1 \right]$$

$$\Rightarrow u_3 = B \left[\cos \left(\sqrt{\frac{n^2 \pi^2}{4L^2}} x_1 \right) - 1 \right]$$

forward this gives:

$$u_3 = B \left[\cos \left(\frac{n\pi}{2L} x_1 \right) - 1 \right]$$

for the case of $n=1$ and P_{cr} :

$$u_3 = B \left[\cos \left(\frac{\pi}{2L} x_1 \right) - 1 \right]$$

with

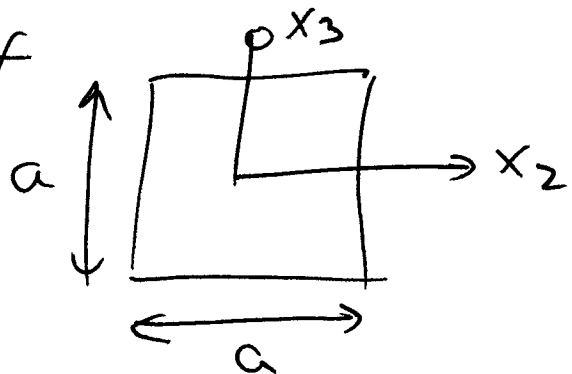
$$P_{cr} = \frac{\pi^2 EI}{4L^2}$$

buckling load

buckling mode

Now look at each of the two cases:

(a) Cross-section of a square



$$A = a^2 \Rightarrow P = \rho a^2$$

$$I = \frac{bh^3}{12} = \frac{a^4}{12}$$

The buckling mode stays the same

$$u_3 = B \left[\cos\left(\frac{\pi}{2L}x_1\right) - 1 \right]$$

but it occurs at a load specific to the configuration:

$$P_{cr} = p_{cr} a^2 = \frac{\pi^2 E \frac{a^4}{12}}{4L^2}$$

critical pressure

$$\Rightarrow p_{cr} = \frac{\pi^2 E a^2}{48L^2}$$

critical pressure for square cross-section

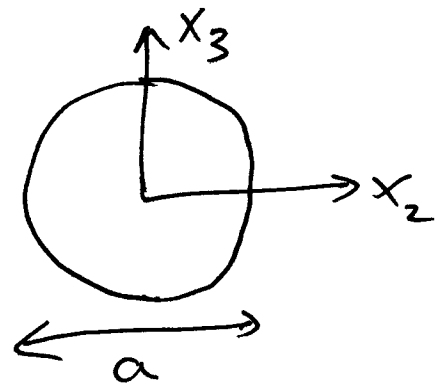
(b) Cross-section of a circle

$$\text{radius} = \frac{a}{2}$$

$$\Rightarrow A = \pi \left(\frac{a}{2}\right)^2 = \frac{\pi a^2}{4}$$

$$\Rightarrow P = p \frac{\pi a^2}{4}$$

$$I = \frac{\pi r^4}{4} \Rightarrow I = \frac{\pi}{4} \left(\frac{a}{2}\right)^4$$



$$\text{giving: } I = \frac{\pi a^4}{64}$$

again, the buckling mode stays the same

$$u_3 = B \left[\cos\left(\frac{\pi}{2L} x_1\right) - 1 \right]$$

but it occurs at a load specific to this cross-section:

$$P_{cr} = P_{cr} \frac{\pi a^2}{4} = \frac{\pi^2 E \frac{\pi a^4}{64}}{4L^2}$$

$$\Rightarrow \boxed{P_{cr} = \frac{\pi^2 E a^2}{64L^2}}$$

critical pressure
for circular
cross-section